

Thermal convection with evolving surface in a rotating icy satellite

Miroslav Kuchta

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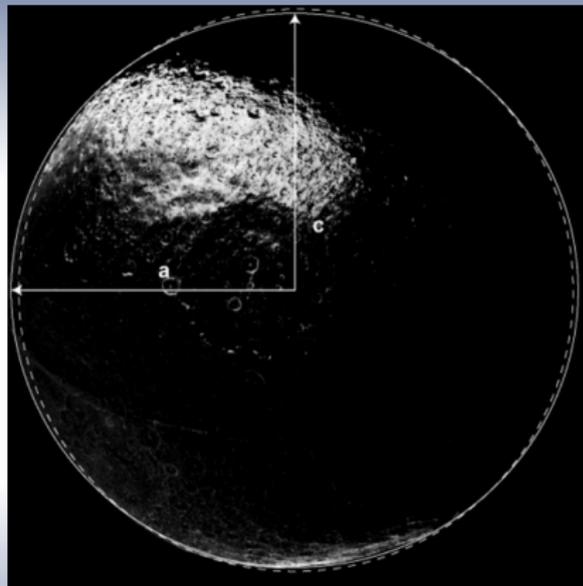
Albedo Dichotomy

- leading dark side - albedo 0.04
- trailing bright side - albedo 0.6
- exogenous origin - volcanism
- endogenous origin - deposits



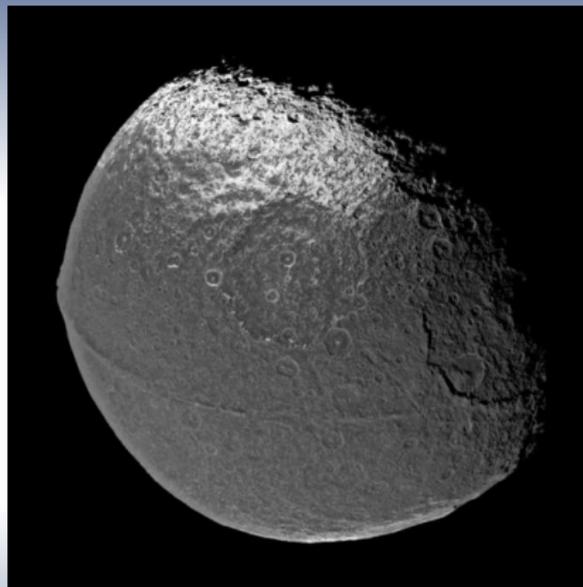
Flattening

- $747.4 \pm 3.1 \times 712.3 \pm 2$ km
consistent with $T_{rot} \approx 16 h$
- observed $T_{rot} = 79.33 d$
- frozen shape: Robuchon et al. [2010], Castillo-Rogez et al. [2007]



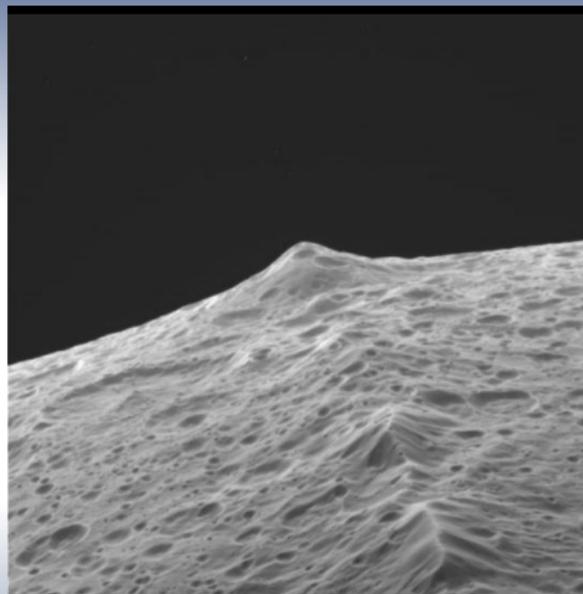
Equatorial Ridge

- endogenous origin: Ip [2006], Levison et al. [2011]
- origin due to contraction: Sandwell and Schubert [2010]
- origin related to convection: Czechowski and Leliwa-Kopystynski [2008]



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Governing Equations

$$\left. \begin{aligned} \nabla \cdot \vec{v} &= 0 \\ -\rho \frac{D\vec{v}}{Dt} - \nabla \pi + \nabla \cdot (\eta (\nabla \vec{v} + (\nabla \vec{v})^T)) + \rho \vec{f} &= 0 \\ \rho_0 c_p \left(\frac{\partial T}{\partial t} + \vec{v} \cdot \nabla T \right) &= \nabla \cdot (k \nabla T) + \mathbf{q} + \boldsymbol{\sigma} : \nabla \vec{v} \end{aligned} \right\} \Omega_t \times (0, T]$$

$$\rho = \rho_0 (1 - \alpha(T - T_0)), \quad \alpha = \text{const} > 0$$

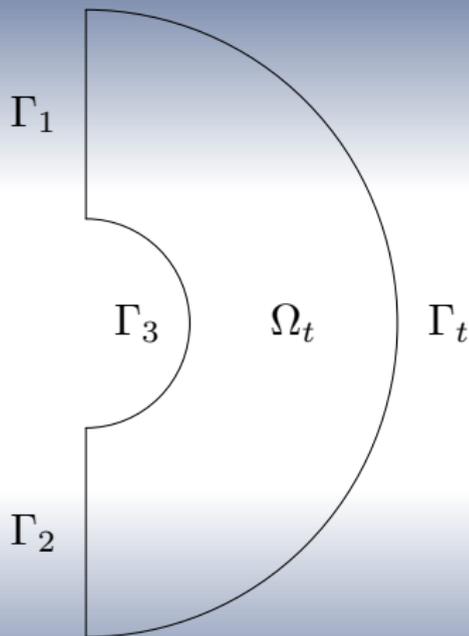
$$\eta = \eta(T, \dot{\epsilon}_{II}, d)$$

$$k = \text{const} > 0, \quad c_p = \text{const} > 0$$

$$\vec{f} = \vec{g}(|\vec{x}|) + \vec{b}(\vec{x})$$

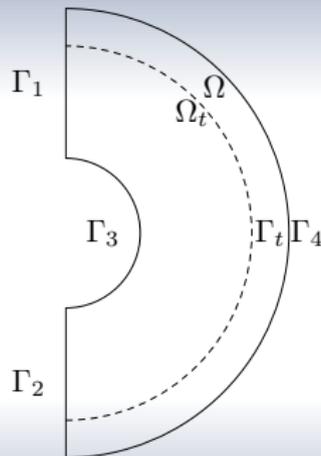
Boundary Conditions

- $(\Gamma_1 \cup \Gamma_2) \times [0, T]$: symmetry
- $\Gamma_3 \times [0, T]$: no-slip, zero heat flux
- $\Gamma_t \times [0, T]$: free-surface, prescribed temperature

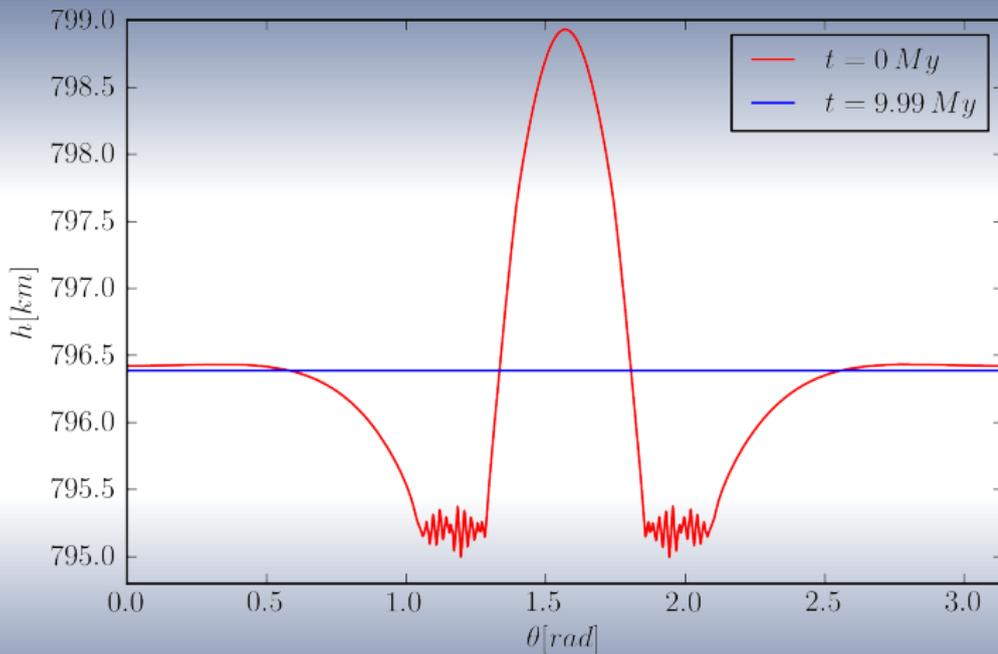


Variable Density Approximation

- $\Omega \supset \Omega_t$, $\partial\Omega = \Gamma_1 \cup \Gamma_2 \cup \Gamma_3 \cup \Gamma_4$
- $\Gamma_4 \times [0, T]$: free-slip, prescribed temperature
- ice := $\{\vec{x} \mid \vec{x} \in \Omega_t\}$
air := $\{\vec{x} \mid \vec{x} \in (\Omega - \Omega_t)\}$
- $\rho_{0,\text{air}} \ll \rho_{0,\text{ice}}$, $\eta_{0,\text{air}} \ll \eta_{0,\text{ice}}$



Surface Tracking

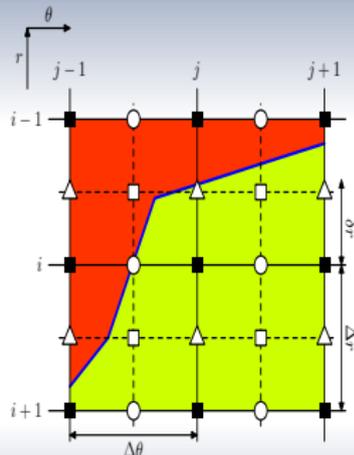


Surface Tracking

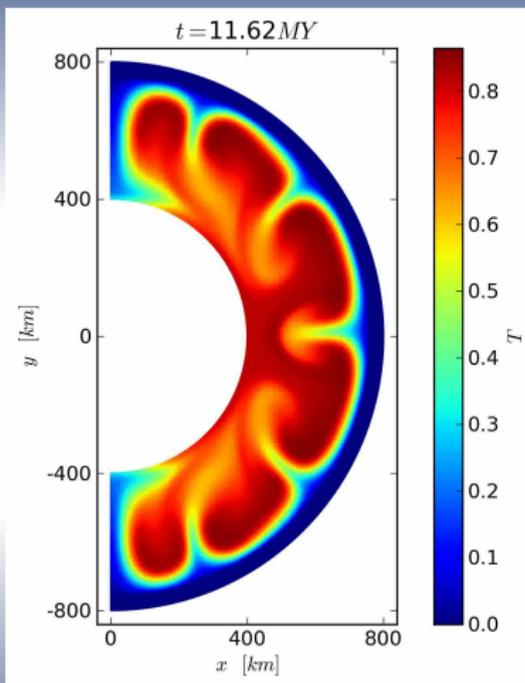
- material surface $H(r, \theta, t) = 0$
- evolution $\frac{\partial H}{\partial t} + \vec{v} \cdot \nabla H + \epsilon \Delta H = 0$
- “nice” surface deformations allow $r = h(\theta, t)$
- symmetry BC on poles

Implementation

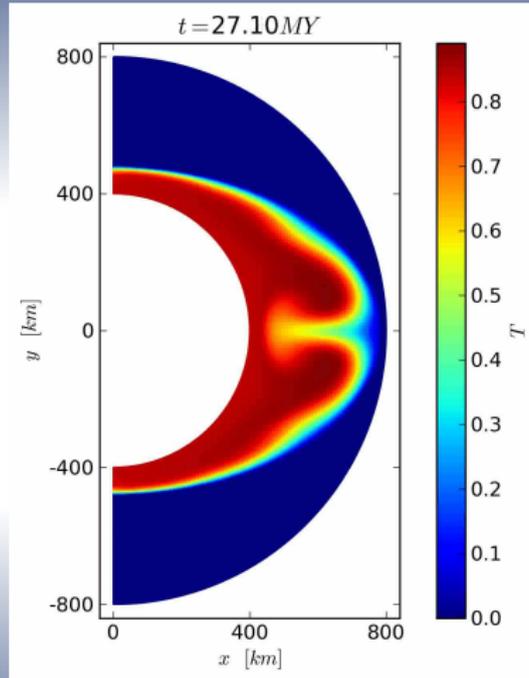
- Stokes problem: solved in Ω , FD on staggered grid
- Balance of energy: solved in Ω_t , FD, semi-implicit with upwinding
- Advection of surface: FD, semi-implicit with upwinding, linear interpolation of ρ near surface



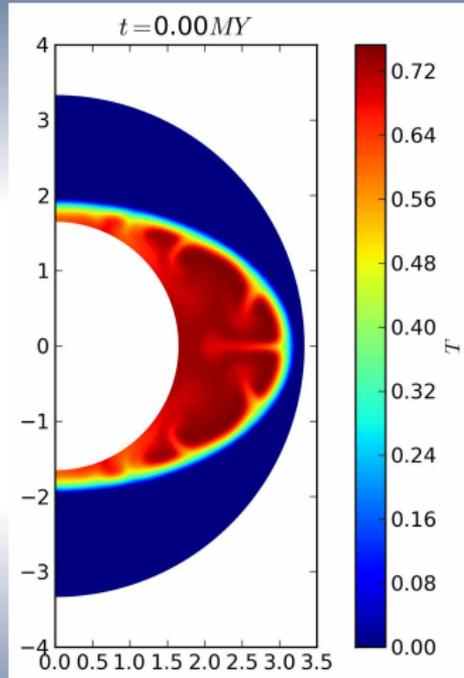
Radiogenic Heating Without Centrifugal Force



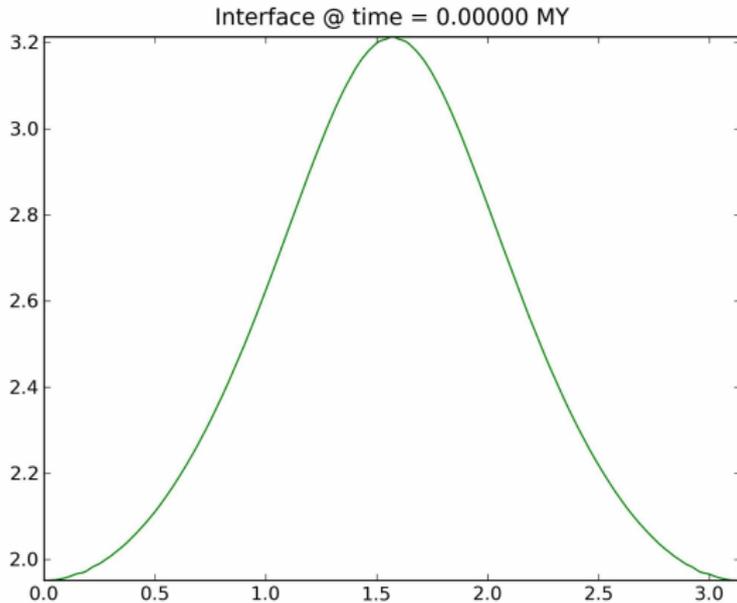
Radiogenic Heating With Centrifugal Force



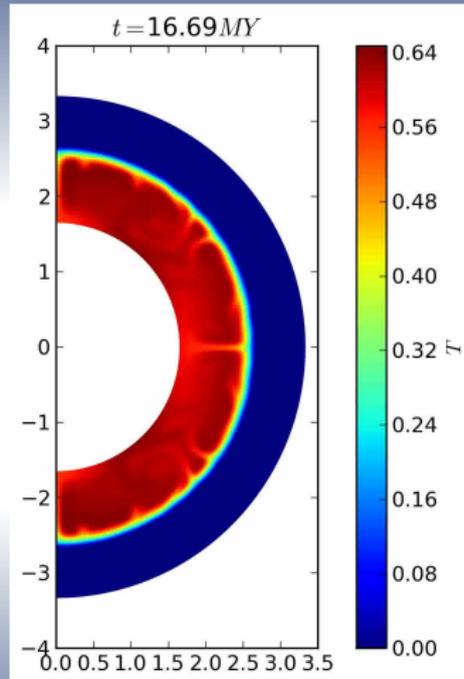
Manual Exponential Despinning



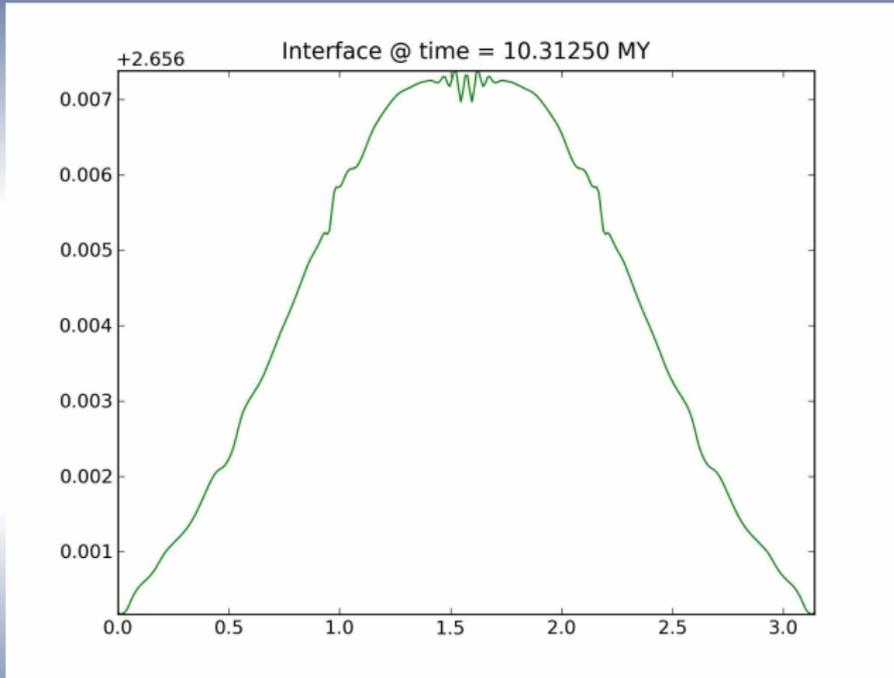
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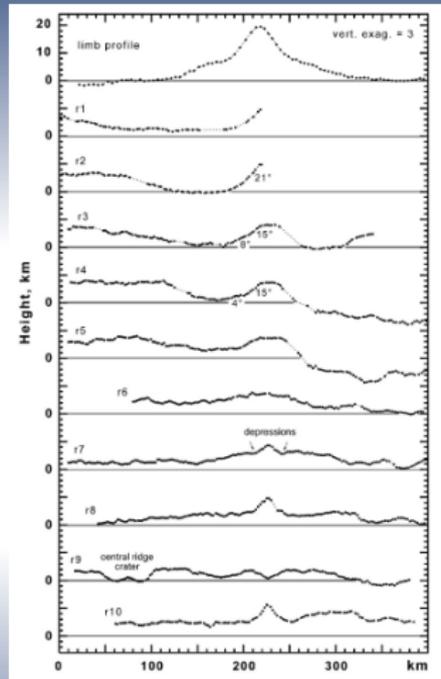
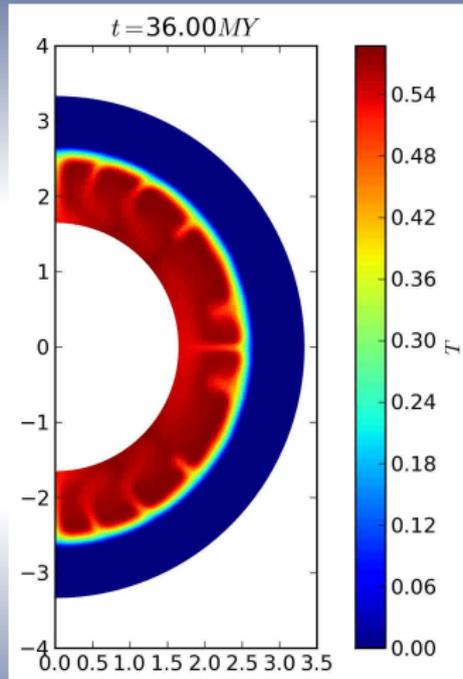
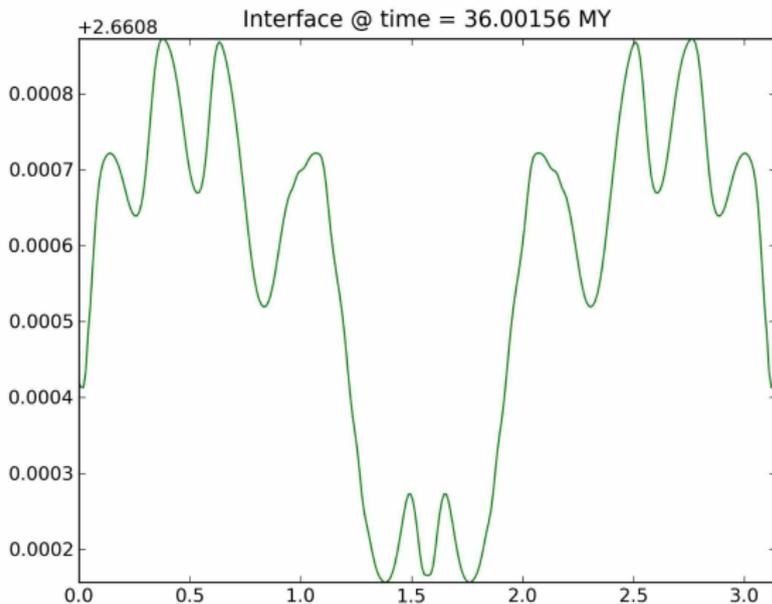


Figure: Iapetian ridge profile Giese et al. [2008]

Manual Exponential Despinning



Manual Exponential Despinning

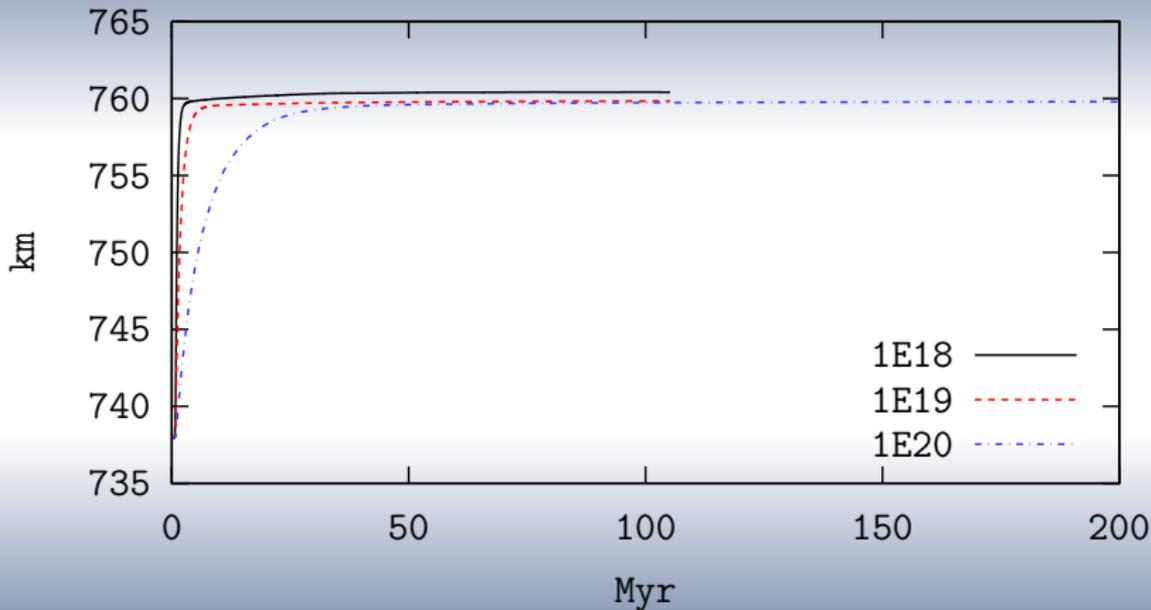


Consistent Despinning

- $E = 50 \text{ kJ.mol}^{-1}$: $\eta_{\text{max}} = 10^{22} \text{ Pa.s}$, $\eta_{\text{air}} = 10^{19} \text{ Pa.s}$

Consistent Despinning

Evolution of Equatorial Radius



Consistent Despinning

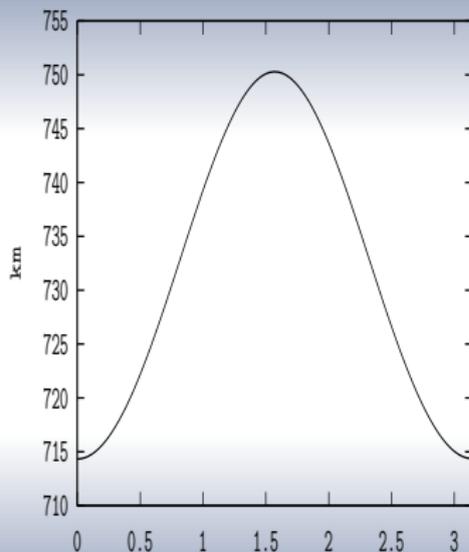
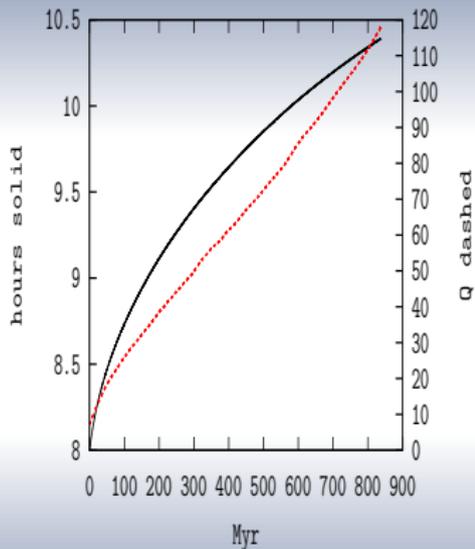
- $E = 50 \times 10^3 \text{ kJ.mol}^{-1}$: $\eta_{\max} = 10^{22} \text{ Pa.s}$, $\eta_{\text{air}} = 10^{19} \text{ Pa.s}$
- IC: $T(\vec{x}, 0) = T_{\text{int}}$ due to SLRI
- spin-rate evolution:

$$\frac{d\omega}{dt} = \frac{3 k_2(t) GM_p^2 a(t)^5}{2 D^6 Q(t) C(t)}$$

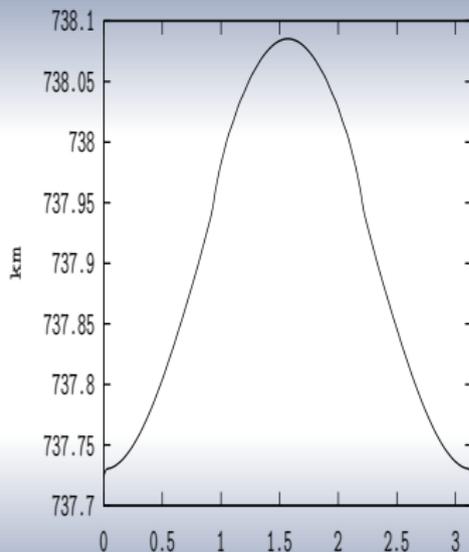
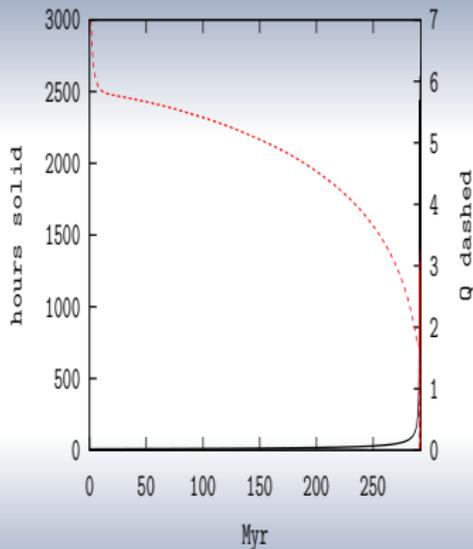
- Andrade rheology:

$$\sigma(t) - \eta_{\text{sym}}(\nabla \vec{v}) = - \int_0^t \frac{\eta}{\mu} \sigma(\tau) d\tau - \int_0^t \alpha \beta (t - \tau)^{\alpha-1} \sigma(\tau) d\tau$$

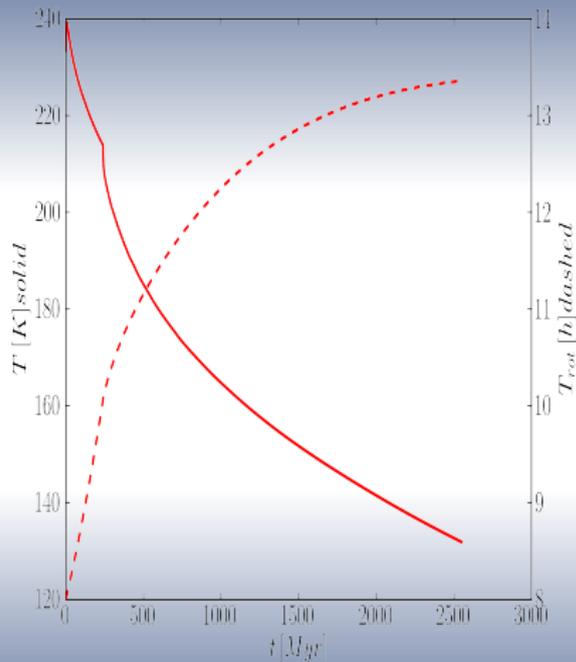
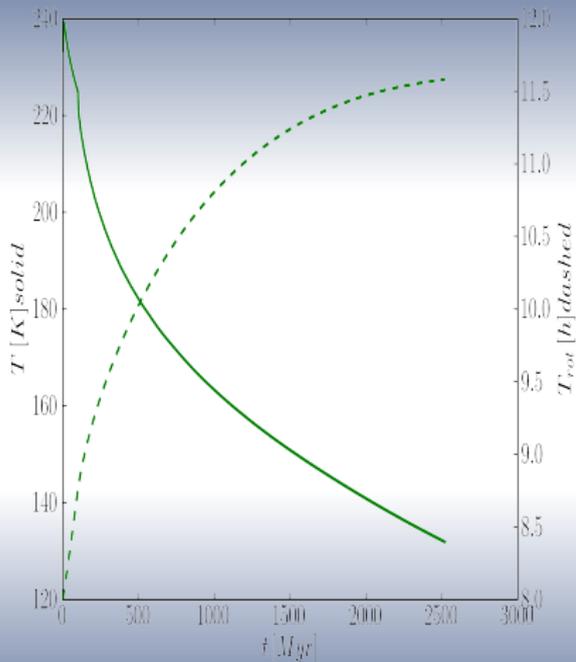
Aarhenius + Plasticity with Convection



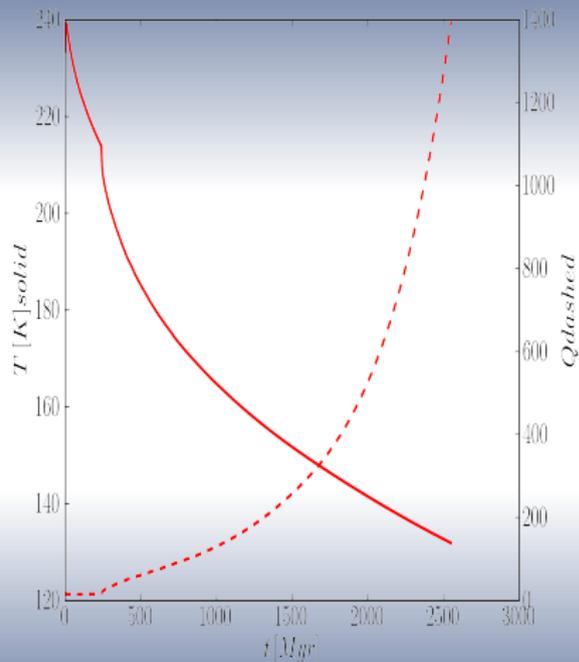
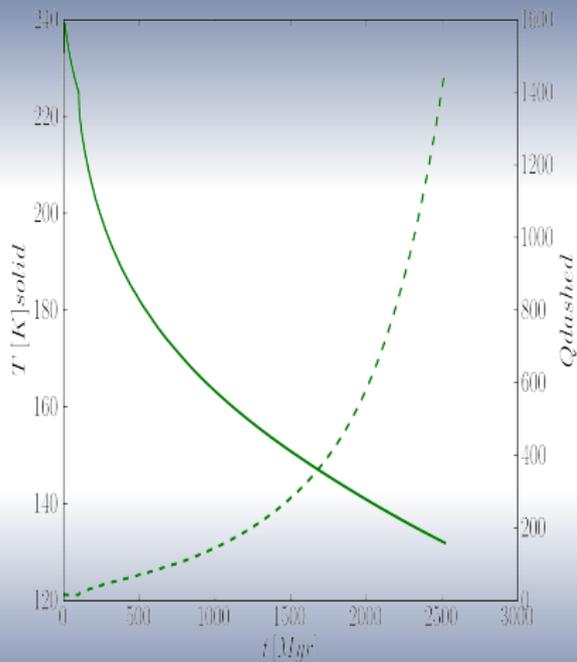
Aarhenius + Plasticity without Convection



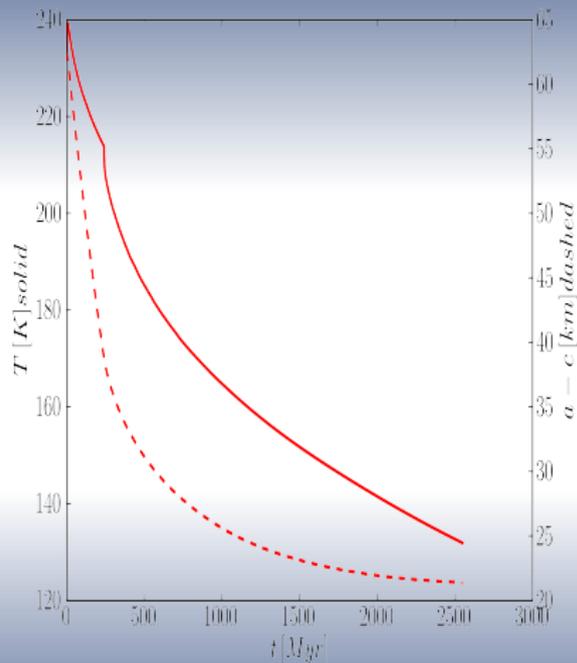
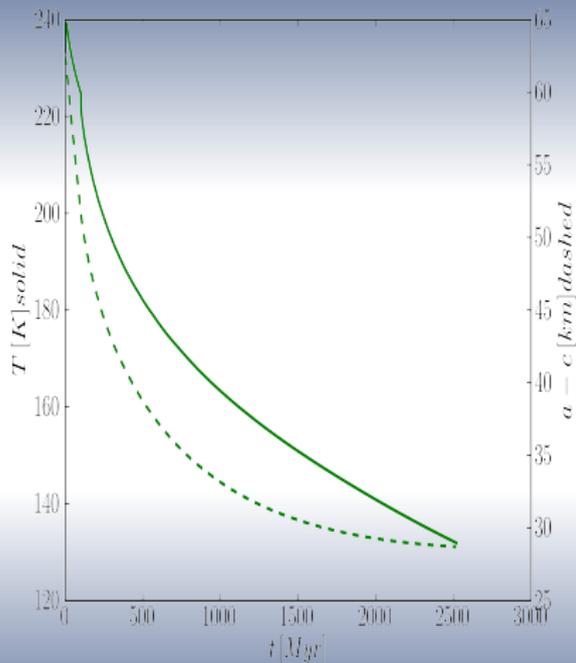
Delayed Convection, $T_{\text{init}} = 240 \text{ K}$



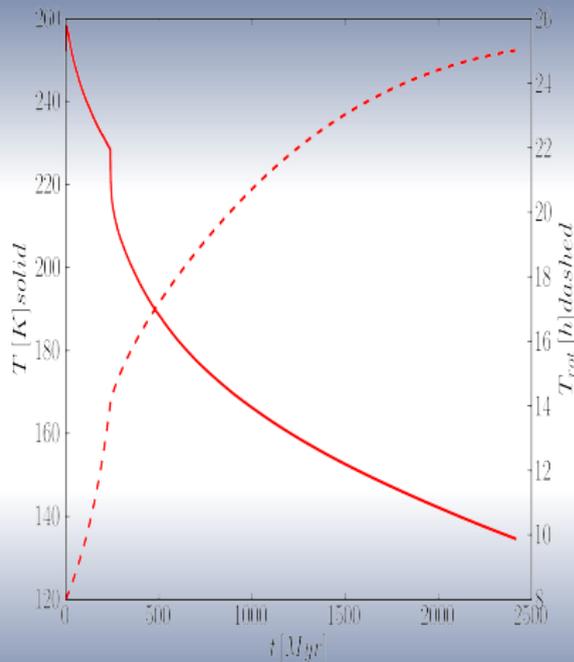
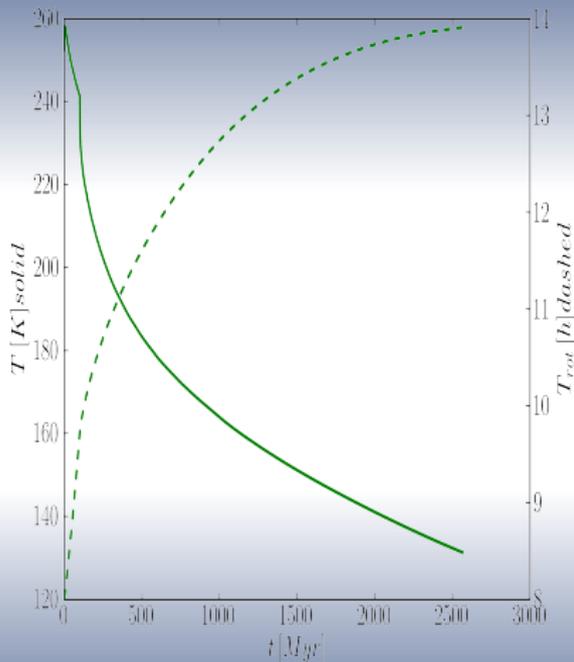
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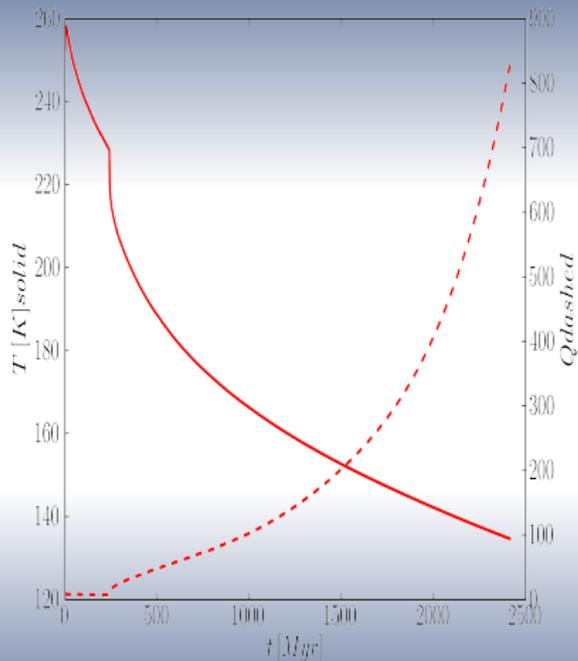
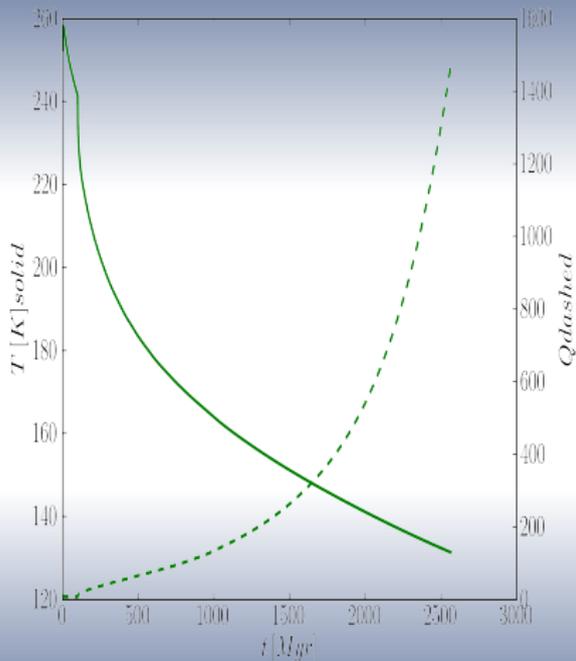
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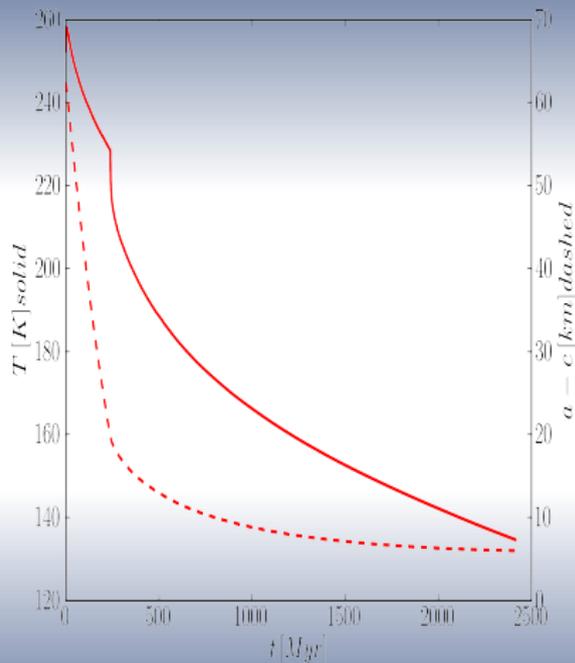
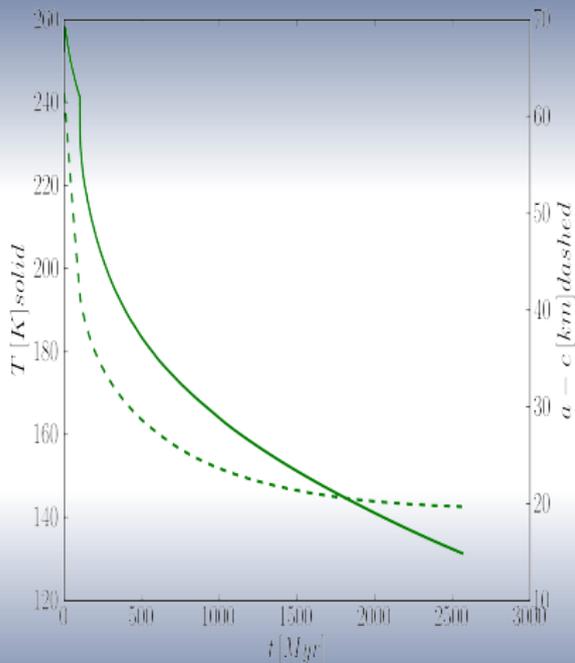
Delayed Convection, $T_{\text{init}} = 260 \text{ K}$



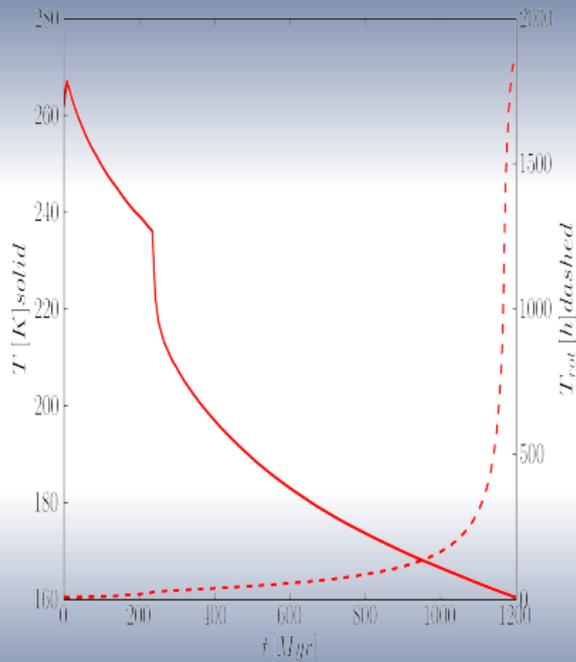
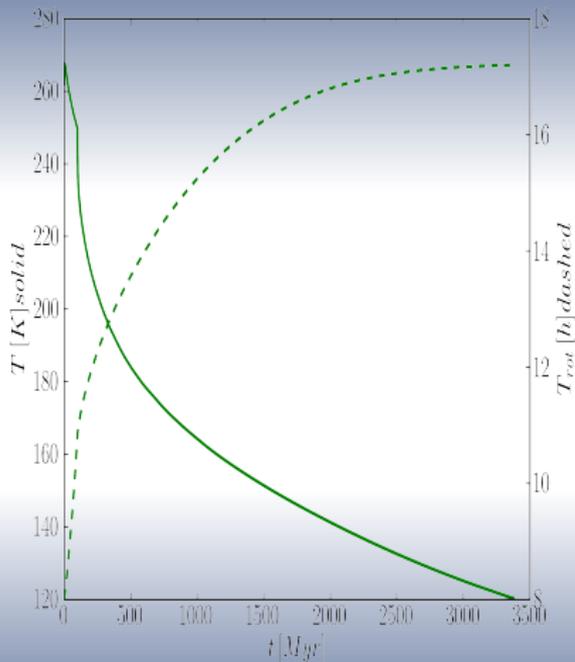
Delayed Convection, $T_{\text{init}} = 260 \text{ K}$



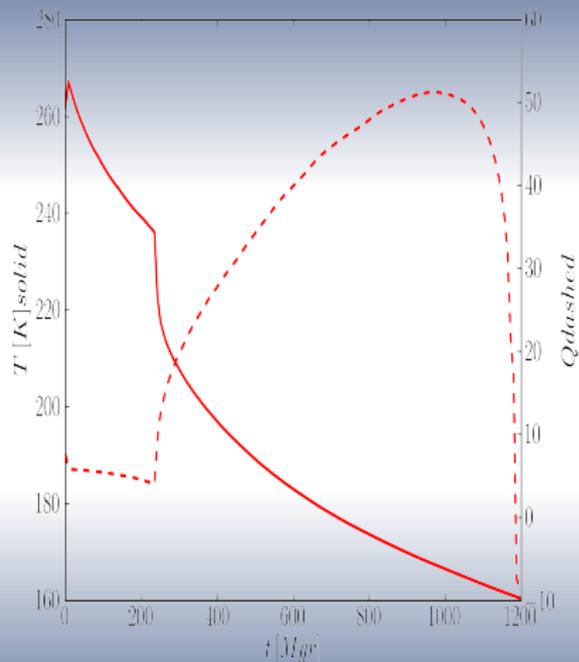
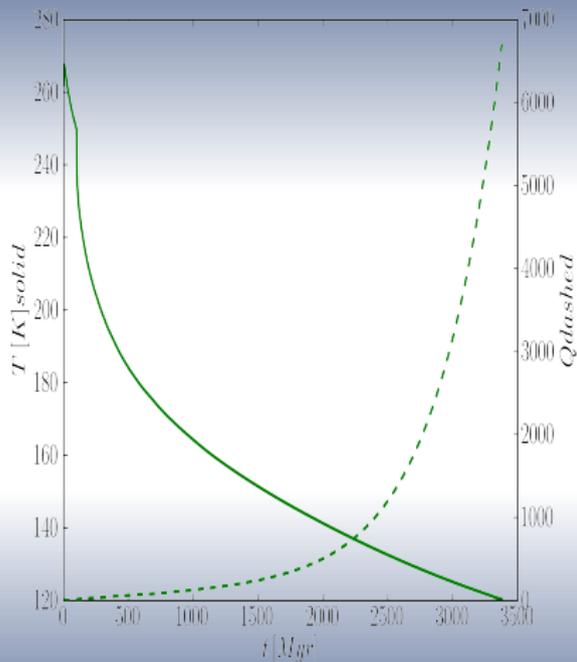
Delayed Convection, $T_{\text{init}} = 260 \text{ K}$



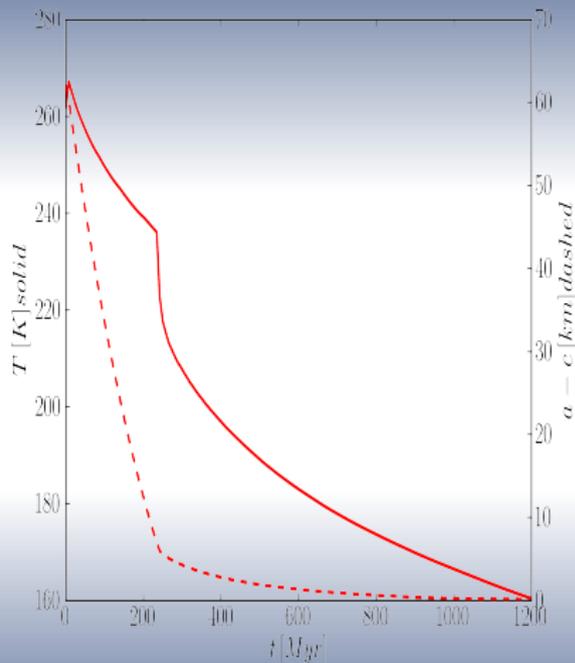
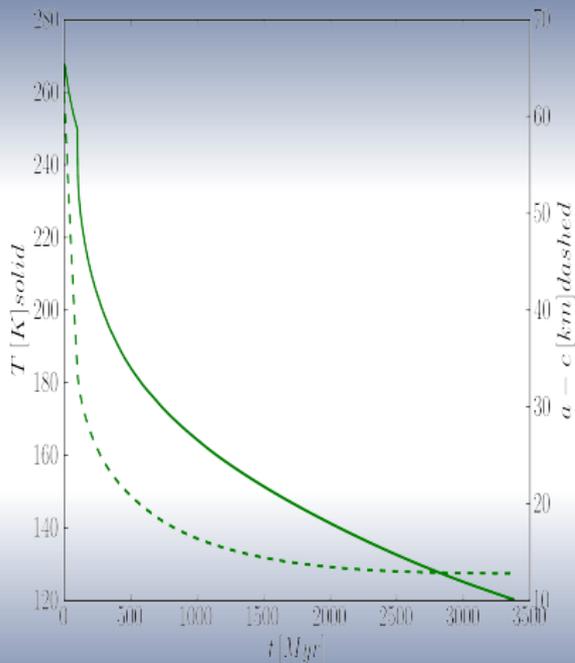
Delayed Convection, $T_{\text{init}} = 270 \text{ K}$



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Deformation Mechanisms of Ice

- dislocation creep, superplastic flow, basal slip creep follow

$$\dot{\epsilon} = A \frac{\sigma^n}{d^p} \exp\left(-\frac{Q + PV}{RT}\right)$$

Regime	A MPa ⁻ⁿ s ⁻¹	n	Q kJ.mol ⁻¹
Dislocation $T < 258$ K	4.0×10^5	4	60
Dislocation $T > 258$ K	6.0×10^{28}	4	180
Superplastic $T > 255$ K	3.9×10^{-3}	1.8	49
Superplastic $T > 255$ K	3.0×10^{26}	1.8	192
Basal slip	5.5×10^7	2.4	60

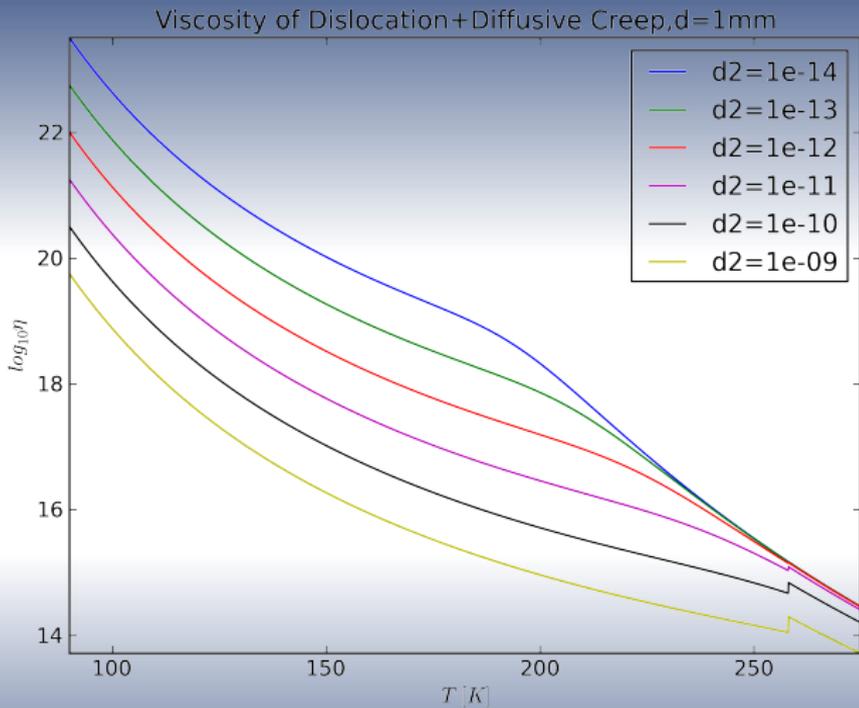
Table: Goldsby and Kehlstedt [2001]

- diffusive flow

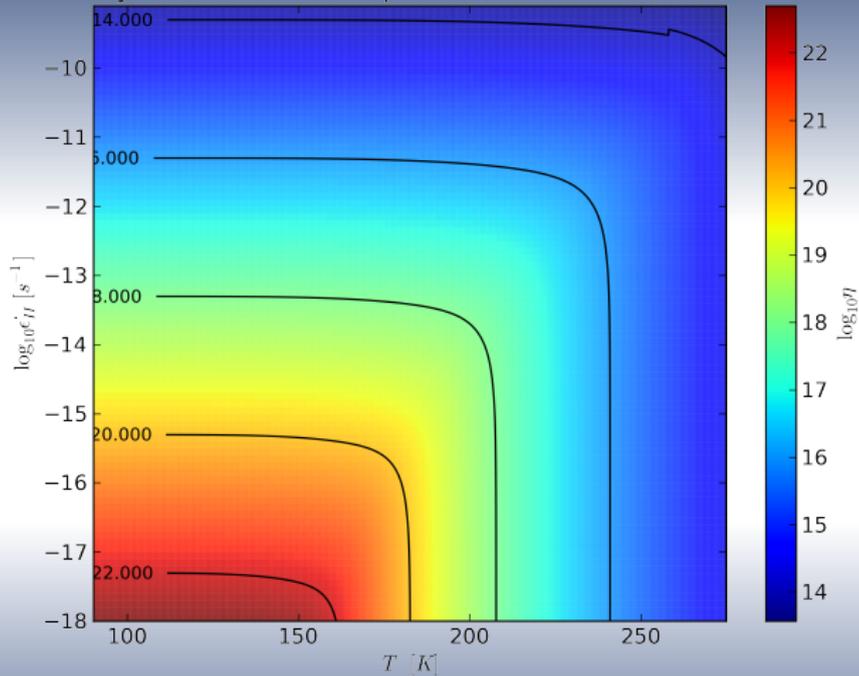
$$\dot{\epsilon} = \frac{42\sigma V_m}{RTd^2} \left(D_{V,0} \exp\left(-\frac{Q_V}{RT}\right) + \frac{\pi\delta}{d} D_{b,0} \exp\left(-\frac{Q_b}{RT}\right) \right)$$

- constitutive equation

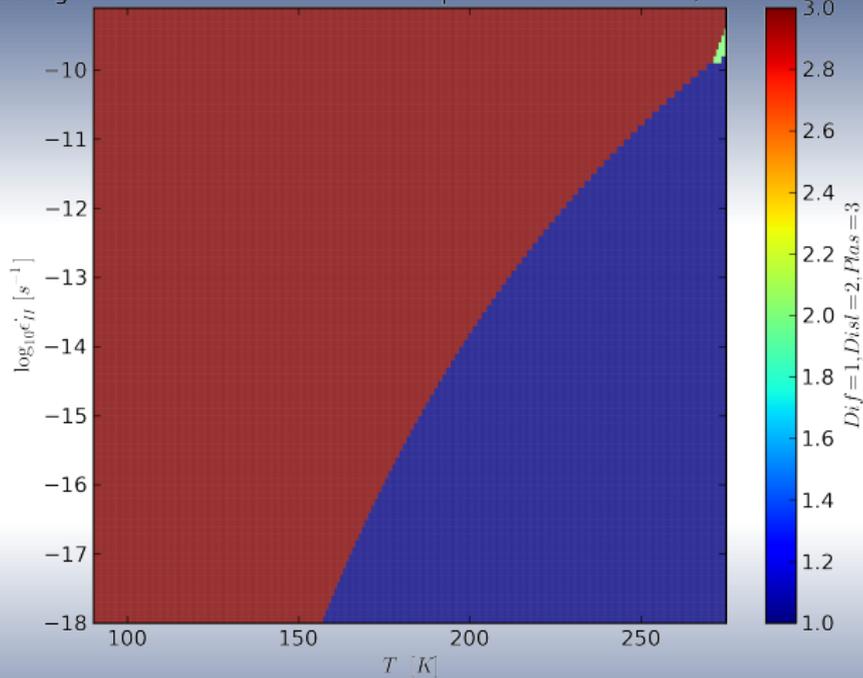
$$\frac{1}{\eta_{TOT}} = \sum_i \frac{1}{\eta_i}$$

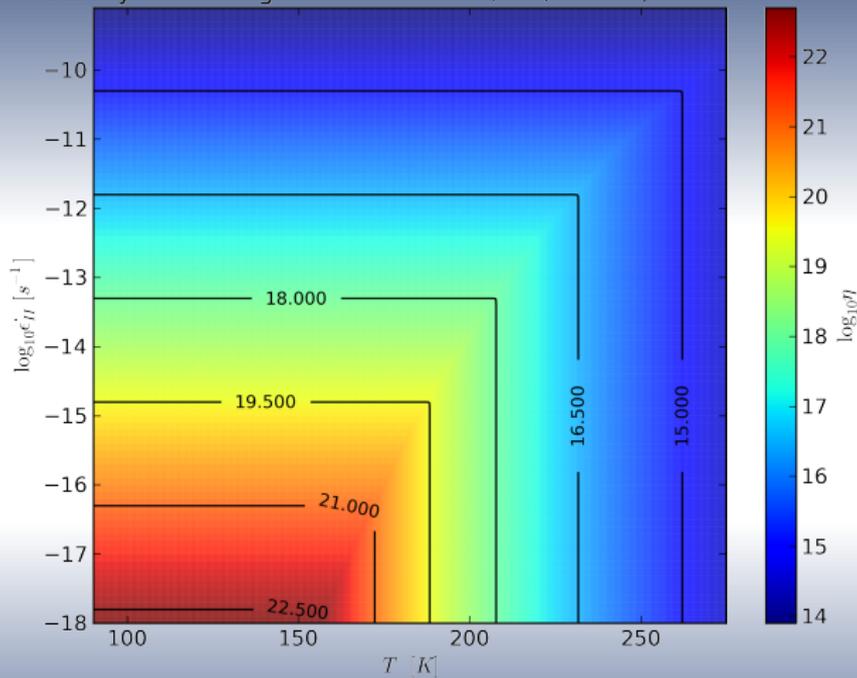


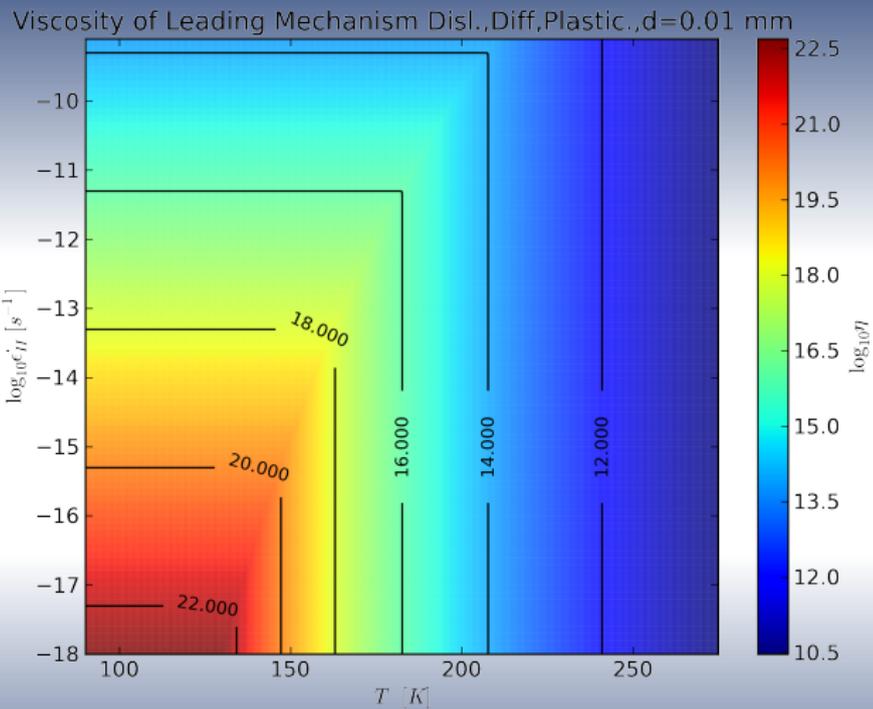
Viscosity of Disl. + Diff. Creep + Plastic. deform., $d=1.00$ mm



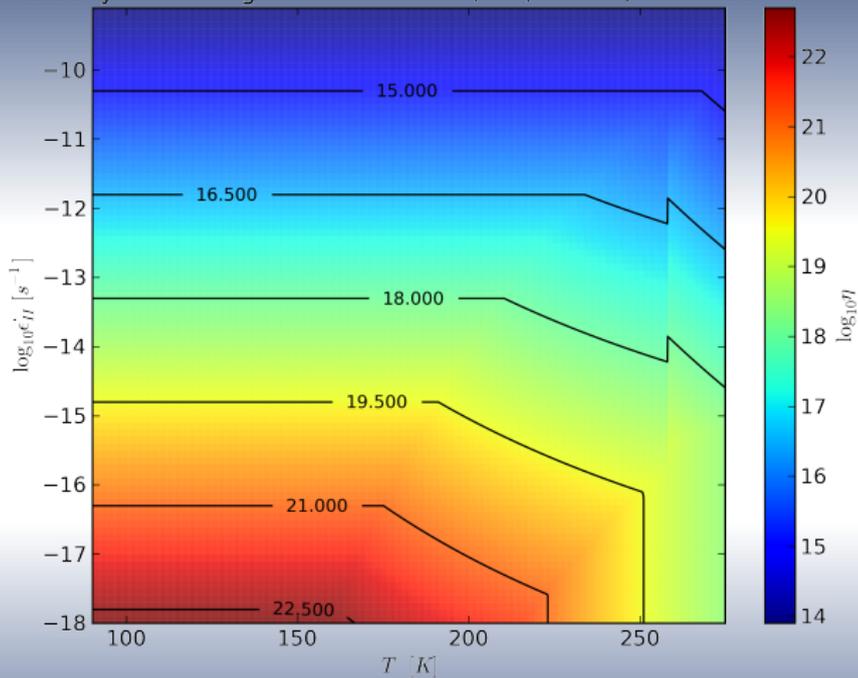
Leading Mechanism Disl. + Diff. Creep + Plastic. deform., $d=1.00$ mm



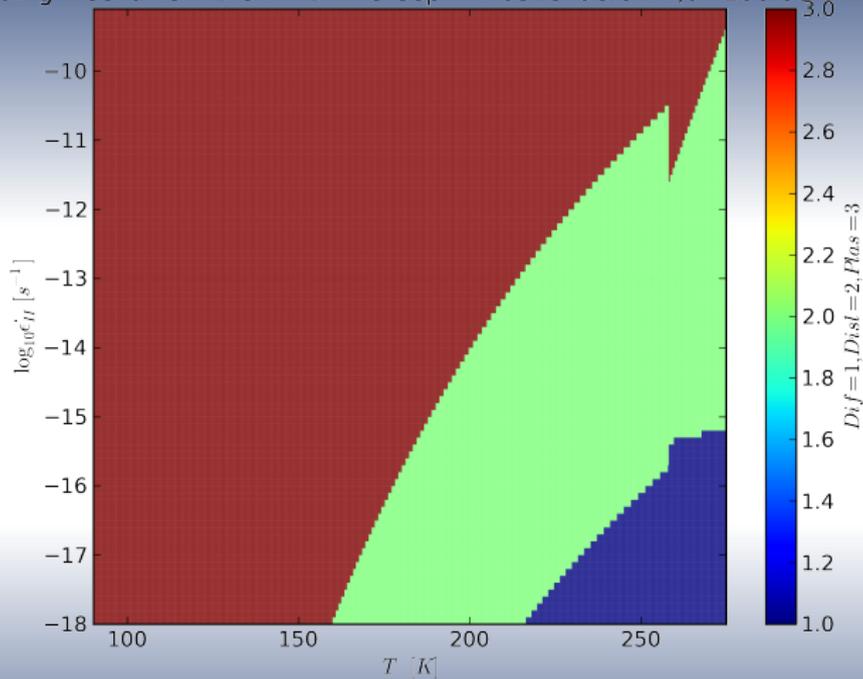
Viscosity of Leading Mechanism Disl., Diff, Plastic., $d=1.00$ mm

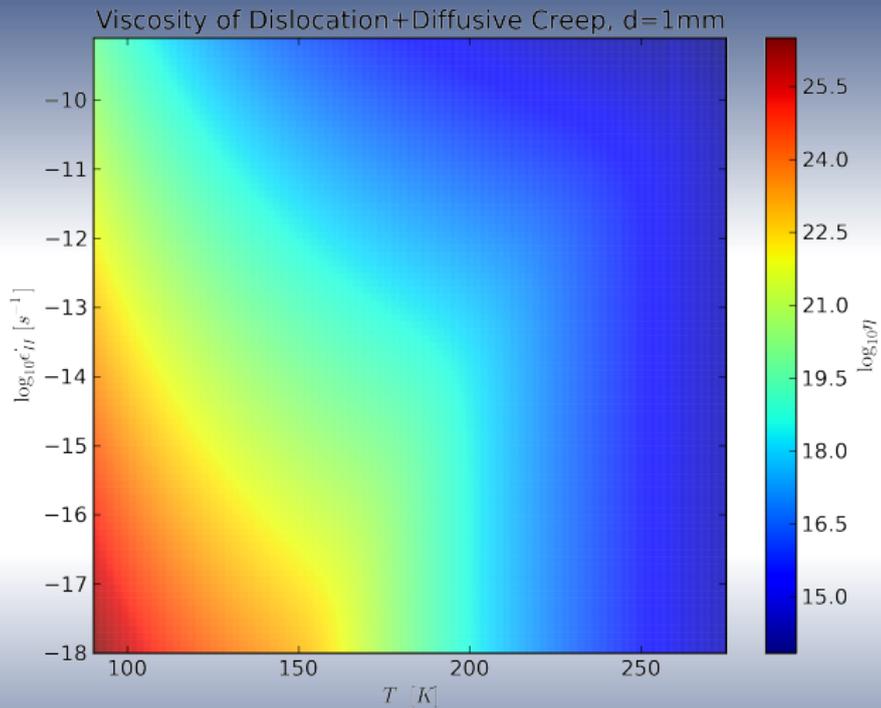


Viscosity of Leading Mechanism Disl., Diff, Plastic., $d=100.00$ mm

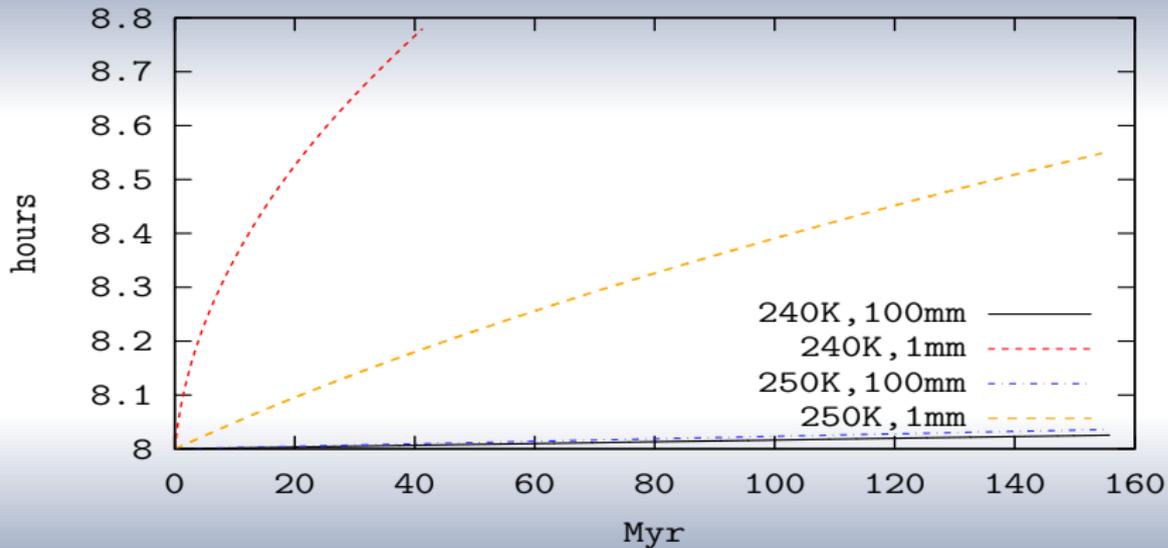


Leading Mechanism Disl. + Diff. Creep + Plastic. deform., $d=100.00$ mm

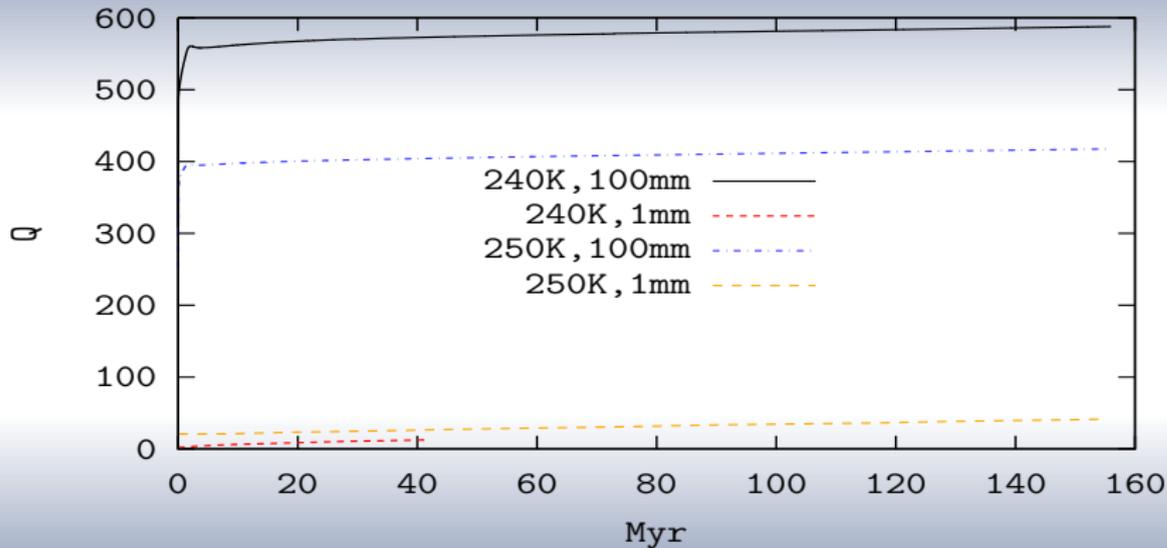




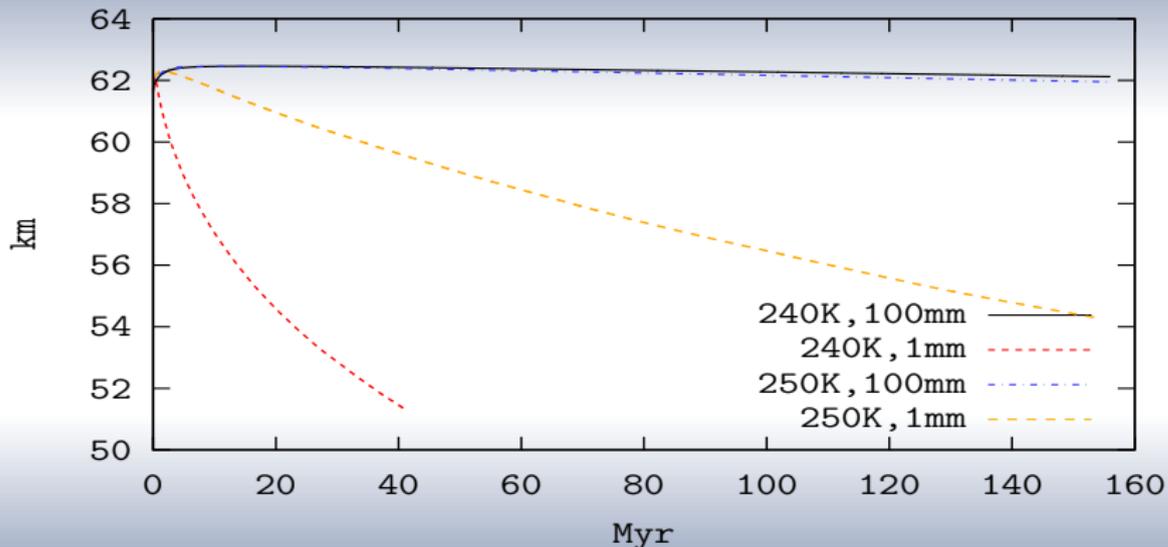
Diffusive Flow + Dislocation Creep + Plasticity



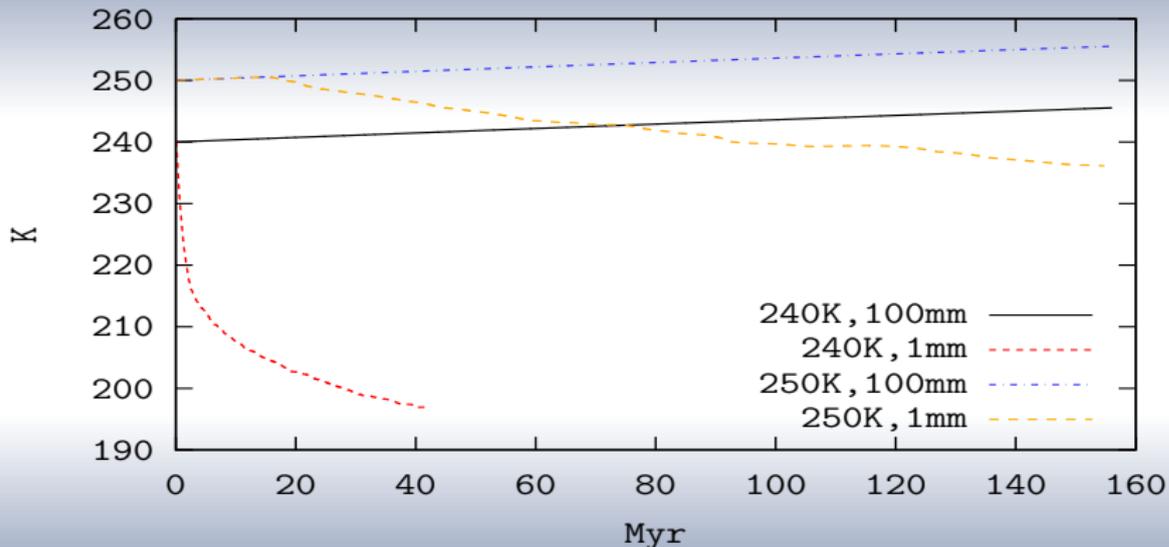
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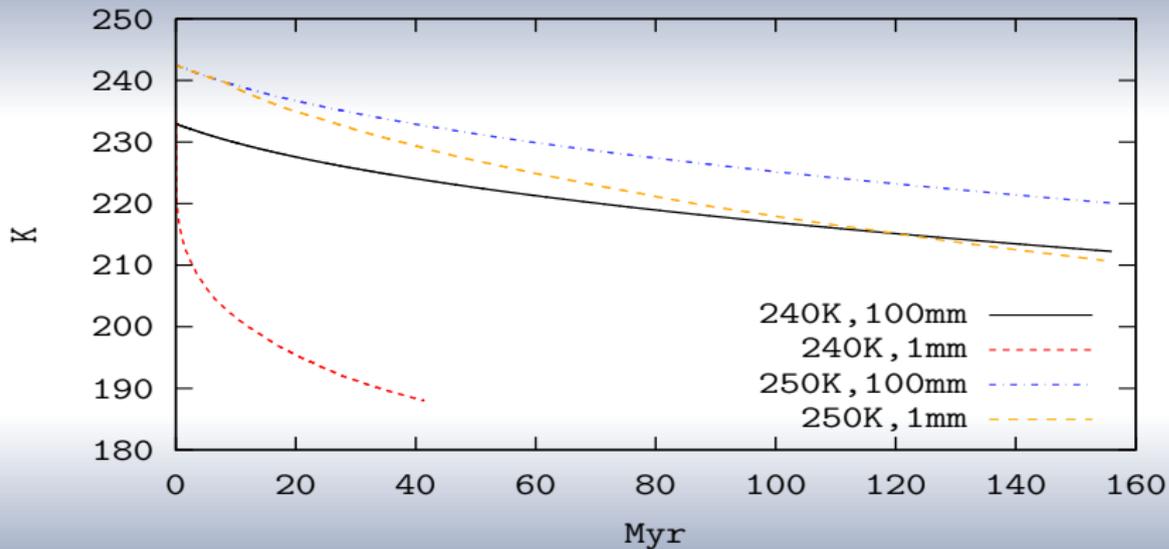
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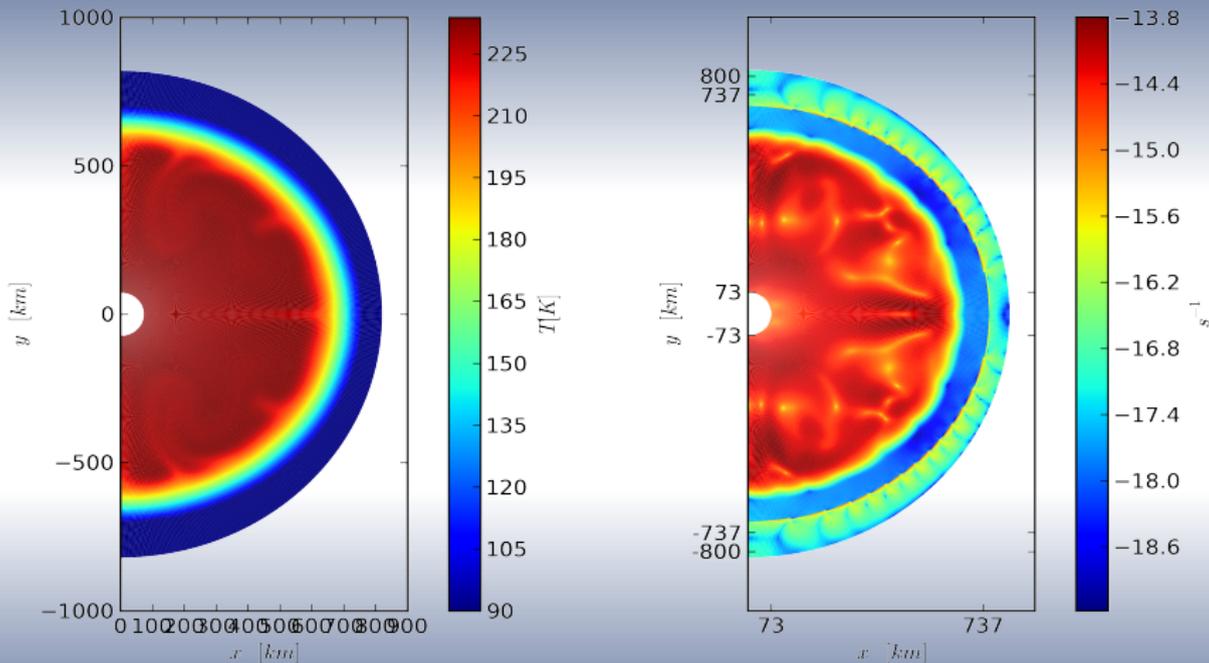
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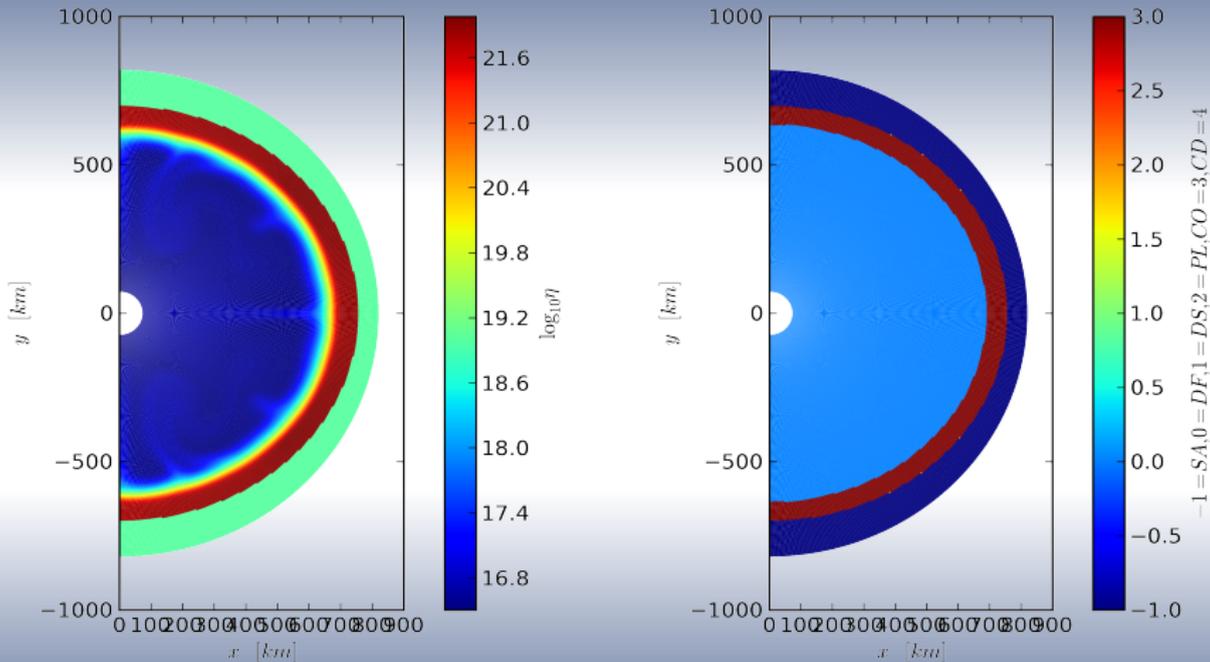
Diffusive Flow + Dislocation Creep + Plasticity



DDP, $d = 1 \text{ mm}$; $T_{\text{init}} = 240 \text{ K}$



DDP, $d = 1 \text{ mm}$; $T_{\text{init}} = 240\text{K}$



- J. C. Castillo-Rogez, D. L. Matson, C. Sotin, T. V. Johnson, J. I. Lunine, and P. C. Thomas. Iapetus' geophysics: Rotation rate, shape, and equatorial ridge. **ICARUS**, 2007.
- L. Czechowski and J. Leliwa-Kopystynski. The Iapetus's ridge: Possible explanations of its origin. **ADVANCES IN SPACE RESEARCH**, 2008.
- Bernd Giese, Tilmann Denk, Gerhard Neukum, Thomas Roatsch, Paul Helfenstein, Peter C. Thomas, Elizabeth P. Turtle, Alfred McEwen, and Carolyn C. Porco. The topography of Iapetus' leading side. **ICARUS**, 193, 2008.
- D. L. Goldsby and D. L. Kehlstedt. Superplastic deformation of ice: Experimental observations. **JOURNAL OF GEOPHYSICAL RESEARCH**, 106, 2001.
- W. H. Ip. On a ring origin of the equatorial ridge of Iapetus. **GEOPHYSICAL RESEARCH LETTERS**, 2006.
- H. F. Levison, K. J. Walsh, A. C. Barr, and L. Dones. Ridge formation and de-spinning of Iapetus via an impact-generated satellite. **ICARUS**, 2011.
- G. Robuchon, G. Choblet, G. Tobie, O. Cadek, C. Sotin, and O. Grasset. Coupling of thermal evolution and despinning of early Iapetus. **ICARUS**, 2010.
- David Sandwell and Gerald Schubert. A contraction model for the flattening and equatorial ridge of Iapetus. **ICARUS**, 2010.